

Series Resistance Both Temperature and Wavelength Dependent in Silicon Solar Cell under Steady State

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Abstract:

The excess minority carrier junction recombination velocity limiting the open circuit S_{foc} and the experimental series resistance of an exposed solar cell to temperature under monochromatic illumination in static regime were determined. It is determined by the photovoltage difference and the open circuit photovoltage V_{oc} . Expression of junction recombination velocity limiting the open circuit has determined the experimental series resistance values. Expressions of photocurrent density and photovoltage are obtained from excess minority carrier density.

Keywords: solar cell - junction recombination velocity - open circuit – wavelength - series resistance - temperature.

I. INTRODUCTION

The series resistance is caused by electrons movement through the emitter and the solar cell base, the contact resistance between the silicon and the metal grids resistance at front and back side [1- 3].

In this paper, we study influence of both temperature and wavelength on series resistance of silicon solar cell in static regime under monochromatic illumination. This resistance and S_{foc} are determined from the I-V characteristic. The experimental series resistance values are also determined from the series resistance curves calibration versus minority carrier junction recombination velocity S_f .

II. THEORY

The studied silicon solar cell is an n^+pp^+ BSF type. It is represented by the following figure:

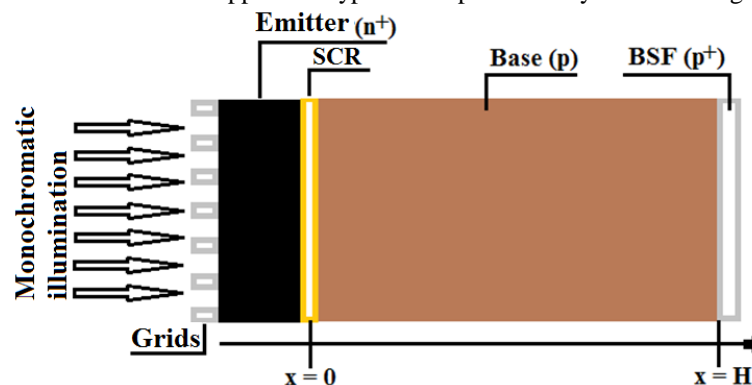


Figure 1 : Silicon solar cell n^+pp^+ type

When the solar cell is illuminated, there is electron-hole pairs generation in the base. The excess minority carrier density in the base is modeled by the following continuity equation:

$$\frac{\partial^2 \delta(x)}{\partial x^2} - \frac{\delta(x)}{L^2} = -\frac{g(x)}{D} \quad (1)$$

With

$\delta(x)$ is the minority carrier density at the depth x in the base

$g(x)$ is the carrier generation rate at the depth x in the base, it is expressed as:

$$g(x) = \phi(\lambda) \alpha(\lambda) (1 - R(\lambda)) e^{-\alpha(\lambda)x} \quad (2)$$

Where $\alpha(\lambda)$ and $R(\lambda)$ represent respectively absorption and reflection coefficient of the material for a given wavelength λ ;

$\phi(\lambda)$ Is the incident photons flux.

The minority carrier diffusion length in the base material is temperature dependent; it is expressed by the following equation.

$$(L(T))^2 = \tau D(T) \tag{3}$$

τ is the minority carrier lifetime.

$D(T)$ is the diffusion coefficient, it's temperature dependent. $D(T)$ can be expressed by Einstein relationship:

$$D(T) = \mu(T) \frac{k_b T}{q} \tag{4}$$

$$\mu(T) \text{ is minority carriers mobility [11], } \mu(T) = 1,43 \cdot 10^9 T^{-2,42} \text{ cm}^2 \text{V}^{-1} \text{s}^{-1} \tag{5}$$

Solution of equation (1) is given by the following terms:

$$\delta(x, \lambda, T) = A \cosh\left(\frac{x}{L(T)}\right) + B \sinh\left(\frac{x}{L(T)}\right) - \frac{\phi(\lambda) \alpha(\lambda) (L(T))^2 (1 - R(\lambda)) e^{(-\alpha(\lambda)x)}}{D(T) [(L(T))^2 (\alpha(\lambda))^2 - 1]} \tag{6}$$

The constants A and B are determined from the boundary conditions:

- At the junction surface, $x = 0$:

$$D(T) \cdot \left. \frac{\partial \delta(x, \lambda, t)}{\partial x} \right|_{x=0} = S_f \cdot \delta(0, \lambda, t) \tag{7}$$

- At the back side surface, $x = H$:

$$D(T) \cdot \left. \frac{\partial \delta(x, \lambda, t)}{\partial x} \right|_{x=H} = -S_b \cdot \delta(H, \lambda, t) \tag{8}$$

S_f is the minority carrier junction recombination velocity, expresses with two components which describe the solar cell operating point and the losses induced by shunt resistance[12,13].

S_b is the minority carrier surface recombination velocity [13], induced by a back surface field.

III. DENSITE DE PHOTOCOURANT

The photocurrent density results from the minority carrier diffusion at the junction and it is given by:

$$J_{ph}(S_f, \lambda, T) = q D(T) \left. \frac{\partial \delta(S_f, \lambda, T)}{\partial x} \right|_{x=0} \tag{9}$$

IV. PHOTOTENSION

The solar cell photovoltage is given by the Boltzmann relation.

$$V_{ph}(S_f, \lambda, T) = V_T \ln \left[\frac{N_b}{(n_i(T))^2} \delta(0, S_f, \lambda, T) + 1 \right] \tag{10}$$

Where V_T is the thermal voltage, it is defined as:

$$V_T = \frac{k_b T}{q} \tag{11}$$

N_b is the doping level and n_i , the intrinsic minority carrier density [14] expressed as follows:

$$n_i(T) = A T^{\frac{3}{2}} \exp\left(-\frac{E_g}{2k_b T}\right) \tag{12}$$

E_g is the energy gap, it's corresponds to the difference between the conduction and valence band energy expressed as: $E_g = E_c - E_v$. $E_g = 1,12 \cdot 1,6 \cdot 10^{-19}$ J

A is a constant. $A = 3,87 \cdot 10^{16} \text{ cm}^{-3} \text{ K}^{-3/2}$

V. I-V CHARACTERISTIC

V.1 Effect of temperature on I-V characteristic

The figure 2 represents the $I(S_f)$ - $V(S_f)$ characteristic of the solar cell under illumination for different temperature values[14].

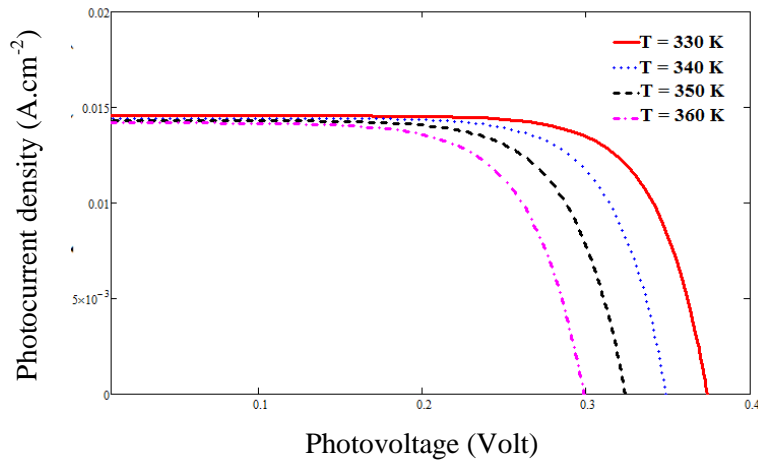


Figure 1: I-V characteristic for different temperature values

We observe that photocurrent density is maximum and constant for low value photovoltage. The maximum photocurrent density is called short-circuit photocurrent density $J_{phsc}(\lambda, T)$. The maximum photovoltage is called open circuit photovoltage $V_{phoc}(\lambda, T)$. For increasing temperature $J_{phsc}(\lambda, T)$ and $V_{phoc}(\lambda, T)$ decreases, slightly for the former and very sensitive for the second one.

V.2 Résistance série

We note in the I-V characteristic that near the open circuit, the voltage across the solar cell is independent of output current. The solar cell is considered as equivalent to a real voltage generator [15, 16]. We have presence of internal resistance called series resistance R_s . R_L is the external load resistance. The electrical equivalent circuit in open circuit situation is then represented [12, 17]:

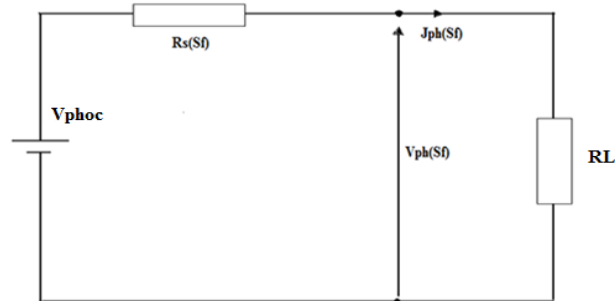


Figure 3: Electrical equivalent circuit in open circuit situation.

By applying the Kirchhoff laws for the previous equivalent circuit, we deduce R_s expression.

$$R_s(S_f, \lambda, T) = \frac{V_{phoc}(\lambda, T) - V_{ph}(S_f, \lambda, T)}{J_{ph}(S_f, \lambda, T)} \quad (13)$$

V.2.1 Temperature effect on series resistance

Figure 4 shows series resistance versus minority carrier junction recombination velocity S_f for different temperature values.

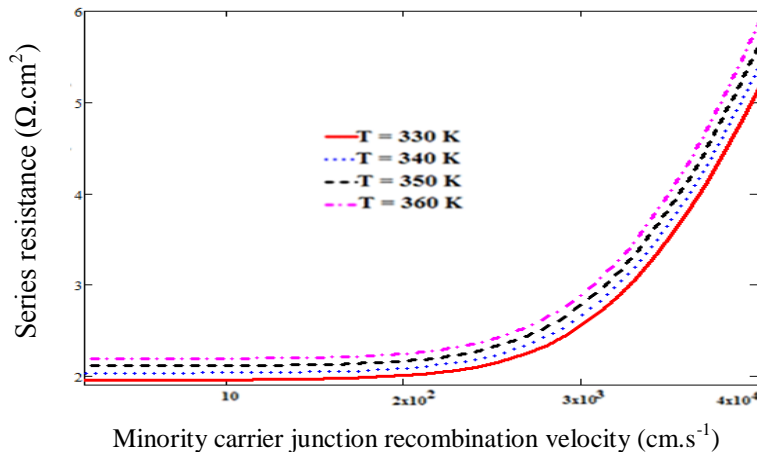


Figure 4: Series resistance versus minority carrier junction recombination velocity S_f for different temperature values.

We observe a bearing series resistance for low Sf values. The series resistance increases exponentially with Sf. And for increasing temperature, the series resistance increases either.

V.2.2 Wavelength effect on series resistance

Figure 5 shows series resistance versus minority carrier junction recombination velocity Sf for different wavelength values.

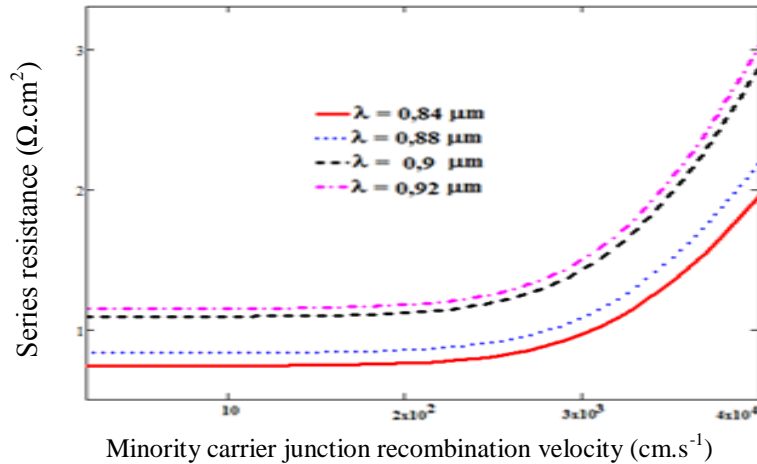


Figure 5: series resistance versus Sf for different wavelength values.

We note that at open circuit situation, series resistance increases gradually with minority carrier junction recombination velocity. We also observe an increase of the series resistance with wavelength: It is explained by the material increasing resistivity.

Figure 6 shows series resistance versus temperature for different wavelength values.

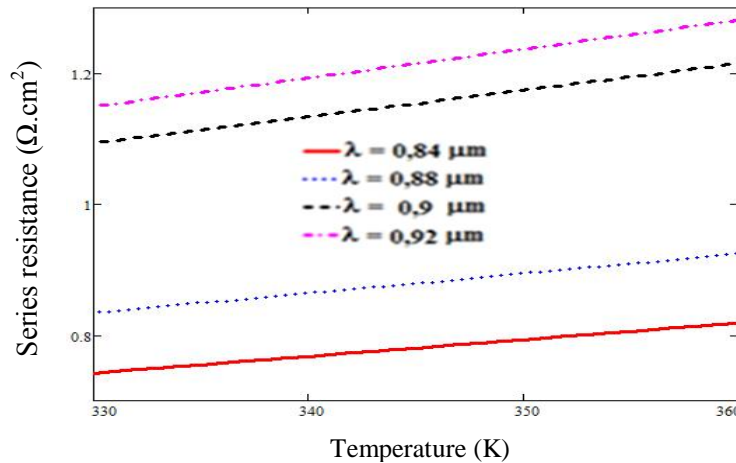


Figure 6: series resistance versus temperature for different high wavelength values (Sf= 10 cm/s)

For increasing temperature, the series resistance increases. It is more important for increasing wavelength values.

Figure 6 allows us to obtain the equation 14 who expressed series resistance versus temperature.

$$Rs_{\lambda}(T) = \beta T + \xi \quad (14)$$

With,

β is the slope and ξ is the intercept and considered as resistance.

β and ξ values are presented for high wavelength by table 1

Table1: β and ξ for high wavelength

$\lambda(\mu\text{m})$	$\beta(\Omega.\text{cm}^2.\text{K}^{-1})$	$\xi(\Omega.\text{cm}^2)$
0,84	$2,55.10^{-3}$	-0,10
0,88	$3,01.10^{-3}$	-0,160
0,9	$4,06.10^{-3}$	-0,247
0,92	$4,22.10^{-3}$	-0,238

V.3 Experimental determination of series resistance

Using the technique of the limiting recombination velocity at the junction Sfoc [18], the expression is derived from the equation 15, at Sf = Sfoc.

$$V_{ph}(S_f, \lambda, T) - V_{phoc}(\lambda, T) = 0 \tag{15}$$

After solving equation 15 we obtain S_{foc} expression given as:

$$S_{foc}(\lambda, T) = \frac{D(T)L(T)[\beta(\lambda, T) + K(\lambda, T)D(T)\alpha(\lambda)]\cos\left(\frac{H}{L(T)}\right) + [D(T)^2\beta(\lambda, T) - L(T)K(\lambda, T)\alpha(\lambda)D(T)]}{D(T)L(T)[K(\lambda, T) - \beta(\lambda, T)]\cos\left(\frac{H}{L(T)}\right) + L(T)[K(\lambda, T) - L(T)\beta(\lambda, T)]\sinh\left(\frac{H}{L(T)}\right)} \tag{16}$$

S_{foc} depend on both wavelength and temperature.

With,

And

$$K(\lambda, T) = \frac{\phi(\lambda)\alpha(\lambda)L(T)^2[1 - R(\lambda)]}{D(T)[L(T)^2\alpha(\lambda)^2 - 1]}$$

$$\beta(\lambda, T) = \frac{n_i(T)^2}{N_b} \left[e^{\left(\frac{V_{co}(\lambda, T)}{V(T)}\right)} - 1 \right] K(\lambda, T)$$

K(λ,T) and β(λ,T) are coefficients, dependent of both wavelength and temperature.

Different temperature values introduced in S_{foc} expression yield corresponding values. The calibration series resistance curve (Figure 4) at S_{foc} values gives either the experimental R_s values who are given by Table 2

Table 2 : S_{foc} and experimental R_s Values with their respective temperature for a given wavelength λ = 0,94 μm.

T(K)	S _{foc} (cm/s)	R _s _{exp} (Ω.cm ²)
330	6,3.10 ³	3
340	6,1.10 ³	3,1
350	5,7.10 ³	3,3
360	5,5.10 ³	3,4

We note in table 2 that for a given wavelength value, S_{foc} decreases for increasing temperature. And for increasing temperature, the experimental series resistance increases.

VI. CONCLUSION

In this study based on temperature effect, we ended up with some results:

- For a given wavelength value, temperature increases causes a reduction of photocurrent density and photovoltage.
- Theoretical and experimental series resistance increases with increasing temperature and for these conditions the Junction surface recombination velocity limiting the open circuit decreases.
- The series resistance increases either with increasing wavelength.

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